Almost-Periodic Solutions of Navier-Stokes Equations and Inequalities

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Introduction

In this talk I would like to present some results, old and new, concerning almostperiodic solutions of Navier-Stokes equations and inequalities, which govern the motion of viscous compressible or incompressible fluids (respectively gases or liquids).

Of the various problems which can be associated with this motion I shall, in what follows, for the sake of simplicity, consider only the one corresponding to a fluid in a bounded 2- or 3-dimensional domain Ω , which boundary Γ constituted by a material surface. Denoting by $\vec{u}(x,t)$ ($x=\{x_1, x_2, x_3\} \in \overline{\Omega}$) the velocity of the fluid, the problem indicated above corresponds, by the limit layer theory, to the homogeneous Dirichlet boundary condition

$$\vec{u}(x,t) = 0 \qquad (x \in \Gamma). \tag{1.1}$$

The following notations will be used in the sequel.

 $\vec{f}(x,t)$ external force acting on the fluid;

p(x,t) pressure;

 $\rho(x,t)$ density; in the incompressible case ($\rho = \text{const}$) I shall assume, for simplicity, $\rho = 1$;

 μ , ζ viscosity coefficients (resp. shear and bulk viscosity);

 \mathscr{D} space of functions (or vectors) $\epsilon C^{\infty}(\overline{\Omega})$ and with compact support in Ω ;

 \mathcal{N} space of vectors $\vec{v} \in \mathcal{D}$ and such that div $\vec{v} = 0$;

 $H^{s} \quad (s \text{ integer} \ge 0) \text{ space of functions (or vectors) square summable in } \Omega, \text{ together} \\ \text{ with their derivatives (in the sense of distributions) of order } \le s;$

 N^s closure of \mathcal{N} in H^s .

The most common mathematical model associated to the motion of a viscous fluid is constituted by the *Navier-Stokes equations* which, in the case of incompressible fluids, take the form

$$\begin{cases} \frac{\partial \vec{u}}{\partial t} - \mu \Delta \vec{u} + (\vec{u}.\text{grad}) \ \vec{u} + \text{grad} \ p = \vec{f} \\ \text{div} \ \vec{u} = 0. \end{cases}$$
(1.2)

while, if the fluid is compressible, are expressed by

$$\begin{cases} \varrho \frac{\partial \vec{u}}{\partial t} + (\zeta + \frac{1}{3}\mu) \text{ grad div } \vec{u} - \mu \Delta \vec{u} + \varrho (\vec{u}.\text{grad}) \vec{u} + \text{grad } p = \varrho \vec{f} \\ \frac{\partial \varrho}{\partial t} + \text{div } (\varrho \vec{u}) = 0 \\ p = p(\varrho). \end{cases}$$
(1.3)

The third equation of (1.3) is an *equation of state* which, in most practical cases, is given by $p = k \varrho^{\gamma} (k, \gamma > 0)$.

It should be noted that (1.2) cannot be considered as a special case of (1.3), since the two systems are essentially different.

Another model associated to viscous incompressible flow corresponds to the *Navier-Stokes inequalities* which are introduced as follows. Observing that the Navier-Stokes equations are non-relativistic and, consequently, do not have any physical meaning when $|\vec{u}|$ approaches the velocity of light, the model (1.2) is equivalent, from a physical point of view, to the one corresponding to the relationships

$$\begin{cases} \frac{\partial \vec{u}}{\partial t} - \mu \Delta \vec{u} + (\vec{u}.\text{grad}) \vec{u} + \text{grad } p = \vec{f} \quad \text{where } |\vec{u}| < c \end{cases}$$

$$(1.4)$$

$$\text{div } \vec{u} = 0 \quad , \quad |\vec{u}| \le c$$

$$\vec{u} = \vec{u} \quad |\vec{u}| \le c$$

 $|\vec{u}|$ continuous at the "interfaces" of the two sets in which resp. $|\vec{u}| < c$ and $|\vec{u}| = c$.

It is well known, on the other hand, from the theory of differential inequalities (see, for instance, Lions [1]) that (1.4) is equivalent to the system

$$\int_{t_1}^{t_2} \int_{\Omega} \left(\frac{\partial \vec{u}}{\partial t} - \mu \Delta \vec{u} + (\vec{u}.\text{grad}) \, \vec{u} + \text{grad} \, p - \vec{f} \right) \, (\vec{u} - \vec{\varphi}) \, dt \, d\Omega \le 0.$$

$$\text{div } \vec{u} = 0$$

$$|\vec{u}| \le c$$
(1.5)

 $V \vec{\varphi}$ such that $|\vec{\varphi}| \leq c$ and $V t_1, t_2 \in (-\infty, +\infty)$.

System (1.5) therefore constitutes an *inequality model* for the problem considered, in the incompressible case. An analogous model could obviously be given for compressible fluids, but it will not be considered here.

In the next section I shall recall some results concerning the almost-periodic solutions of the three models presented; it is however useful to first briefly summarize the main existence and uniqueness theorems of a solution of (1.2), (1.3), (1.5) satisfying (1.1)and the initial conditions

$$\vec{u}(x, 0) = \vec{u}_0(x)$$
(incompressible case)
$$\vec{u}(x, 0) = \vec{u}_0(x), \ \varrho(x, 0) = \varrho_0(x)$$
(compressible case)

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These theorems represent, in fact, the first step in the study of the almost-periodic solutions.

The solutions will always be intended in the sense of distributions, while I shall not, for the sake of simplicity, indicated explicitly the functional spaces in which the solutions are found, or the assumptions on the data.

Considering system (1.2), Hopf [2] proved the global (in time) existence of a solution in any space dimension; the uniqueness of such a solution can however be guaranteed only in 2 dimensions (Lions and Prodi [3]). An existence and uniqueness theorem in Ω × (0, *T*), Ω 3-dimensional, holds provided \vec{f} is "sufficiently small" (Kieselev and Ladyzenskaja [4]).

One can, on the other hand, prove a global existence and uniqueness theorem for the solution in $\Omega \times (0, T)$ of (1.5) (Prouse [5]).

In the compressible case, only a local existence and uniqueness theorem holds (Valli [6]); in order to obtain global existence and uniqueness, one must assume that \vec{f} is "sufficiently small" (Marcati and Valli [7]).

Almost-periodicity theorems

The models introduced in the preceding section all correspond to dissipative problems, and the study of their almost-periodic solutions follows therefore from the guidelines given, for ordinary dissipative differential equations, by Favard [8] and Amerio [9] respectively in the linear and non linear case.

In the theory of almost-periodic solutions of partial differential equations, vector valued functions play an essential role, together with the concepts of weakly almost-periodic and S^{p} -Stepanov almost-periodic functions. For these concepts and for the basic definitions and properties of functions with values in a Banach space, I refer to the note by L. Amerio which appears in the present volume (see also Amerio, Prouse [10]).

While the details of the proofs of the existence and uniqueness of an almost-periodic solution, under the assumption that $\vec{f}(t)$ is almost-periodic, are obviously different for the three models considered, the basic scheme is similar and consists essentially of the following steps:

- a) A global existence theorem in $[t_0, +\infty)$;
- b) An existence and uniqueness theorem of a solution $\tilde{\vec{u}}(t)$ (or $\{\tilde{\vec{u}}(t), \tilde{\varrho}(t)\}$) bounded on $J = (-\infty, +\infty)$ (assuming $\vec{f}(t)$ bounded on J);
- c) The proof that $\vec{\tilde{u}}(t)$ ({ $\vec{\tilde{u}}(t)$, $\tilde{\varrho}(t)$ }) is weakly almost-periodic if $\vec{f}(t)$ is weakly almost-periodic;
- d) The proof that the range of $\vec{\tilde{u}}(t)$ ($\{\vec{\tilde{u}}(t), \tilde{\varrho}(t)\}$) is relatively compact if $\vec{f}(t)$ is almost-periodic.

Observe that point a) corresponds essentially to the results recalled in the preceding section, setting $T = +\infty$.

Assuming that f(t) is S²-Stepanov almost-periodic, the following theorems then hold.

THEOREM I (Prouse [11]): If Ω is 2-dimensional, $\vec{f}(t) \in L^{\infty}(J; L^2)$ and is "sufficiently small", (1.1), (1.2) admit a unique solution $\vec{u}(t)$ which is N^0 -Bohr and N^1 -S²-Stepanov almost-periodic.

THEOREM II (Foias [12], Heywood [13]): If Ω is 3-dimensional and of class C^3 , $\tilde{f}(t) \in L^2_{loc}(J, N^1) \cap H^1_{loc}(J, (N^1)^*)$ and is "sufficiently small", then (1.1), (1.2) admit a unique solution $\tilde{u}(t)$ which is N^0 -Bohr and N^1 -S²-Stepanov almost-periodic.

THEOREM III (Marcati and Valli [7]): If Ω is 3-dimensional and of class C^4 , $p \in C^3$, p' > 0, $\vec{f}(t) \in L^2_{loc}(J; H^1) \cap H^1_{loc}(J; H^{-1})$ and is "sufficiently small", then (1.1), (1.3) admit a unique solution $\{\vec{u}(t), \vec{\varrho}(t)\}$ with $\vec{u}(t) H^1$ -Bohr and H^2 -S²-Stepanov almost-periodic, $\vec{\varrho}(t) L^2$ -Bohr and H^2 -S²-Stepanov almost-periodic.

THEOREM IV (Prouse [14]): If Ω is 3-dimensional, $\tilde{f}(t) \in L^{\infty}(J; L^2)$ and is "sufficiently small", then (1.1), (1.4) admit a unique solution $\tilde{u}(t)$ which is N^0 -Bohr and N^1 -S²-Stepanov almost-periodic.

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